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Buckling analysis of isotropic and orthotropic plates using a novel four variable refined plate theory

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Abstract. The current research presents a buckling analysis of isotropic and orthotropic plates by proposing a new four variable refined plate theory. Contrary to the existing higher order shear deformation theories (HSDT) and the first shear deformation theory (FSDT), the proposed model uses a new displacement field which incorporates undetermined integral terms and involves only four variables. The governing equations for buckling analysis are deduced by utilizing the principle of virtual works. The analytical solution of a simply supported rectangular plate under the axial loading has been determined via the Navier method. Numerical investigations are performed by using the proposed model and the obtained results are compared with CPT solutions, FSDT solutions, and the existing exact solutions in the literature. It can be concluded that the developed four variable refined plate theory, which does not use shear correction coefficient, is not only simple but also comparable to the FSDT.

Keywords: refined plate theory; buckling analysis; isotropic plate; orthotropic plate

1. Introduction

The stability of rectangular plates has been a topic of investigation in engineering structures for more than a century (Matsunaga 2009, Bachir Bouiadjra *et al.* 2013, Altunsaray and Bayer 2014, Afsharmanesh *et al.* 2014, Bakora and Tounsi 2015, Bouguenina *et al.* 2015, Nguyen *et al.* 2015, Tagrara *et al.* 2015, Tebboune *et al.* 2015, Larbi Chaht *et al.* 2015, Musa 2016, Rajanna *et al.* 2016, Bouderba *et al.* 2016, Yousefitabar and Matapouri 2016, Chikh *et al.* 2016, Eltaher *et al.* 2016). A great number of exact solutions for isotropic and orthotropic plates have been proposed, and the most known are documented in Timoshenko and Woinowsky-Krieger (1959), Timoshenko and Gere (1961), Bank and Yin (1996), Kang and Leissa (2005), Hwang and Lee (2006) and Musa (2016). In company with investigations of stability response of plate, a great number of plate models have been proposed. The simplest one is the classical plate theory (CPT) which ignores the transverse shear influences. This model is not suitable for the thick and orthotropic plate with

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important degree of modulus ratio. To avoid this limitation, the shear deformation theory which considers the transverse shear influences is introduced by several researchers. Indeed, the FSDT (first shear deformation theory) and HSDT (high shear deformation theory) were proposed as improvement of the CPT. The FSDT is based on Reissner (1945) and Mindlin (1951) and consider the transverse shear influences by the way of linear distribution of the displacements across the thickness. Many investigations have been presented in different scientific articles by employing FSDT for the free vibration behavior of composite plates (Yan et al. 1966, Whitney 1969, Whitney and Pagano 1970, Ambartsumyan 1970, Sun and Whitney 1973, Bert and Chen 1978, Reddy 1979, Noor and Burton 1989a, b, Kant and Swaminathan 2001). Kant and Swaminathan (2001) extended the FSDT presented by Whitney and Pagano (1970) for the dynamic response of laminated composite and sandwich plates. Sadoune et al. (2014) developed a novel FSDT for mechanical behavior of laminated plates. Meksi et al. (2015) proposed a new simple FSDT with only four variables for static and vibration analysis of functionally graded plates. Bellifa et al. (2016) used a new FSDT for bending and dynamic analysis of functionally graded plates. Hadji et al. (2016) analyzed a FG beam using a new first-order shear deformation theory. Yaghoobi et al. (2014) proposed an analytical study on post-buckling and nonlinear free vibration analysis of FG beams resting on nonlinear elastic foundation under thermo-mechanical loadings using VIM. Bourada et al. (2012) presented a new four-variable refined plate theory for thermal buckling analysis of FG sandwich plates. It is noted that various HSDTs are also used (Mahapatra et al. 2016a, b, Katariya and Panda 2016, Mehar and Panda 2016, Panda and Katariya 2015, Kar and Panda 2014, Panda and Singh 2013a, b, c, 2010a, b, 2011) to investigate the geometrical deformation of composite structures under the large geometrical distortion. Tounsi and his co-workers (Benachour et al. 2011, Houari et al. 2013, Tounsi et al. 2013, Ould Larbi et al. 2013, Saidi et al. 2013, Bousahla et al. 2014, Draiche et al. 2014, Fekrar et al. 2014, Hadji et al. 2014, Khalfi et al. 2014, Zidi et al. 2014, Klouche Djedid et al. 2014, Ait Yahia et al. 2015, Zemri et al. 2015, Sallai et al. 2015, Bennai et al. 2015, Bouchafa et al. 2015, Meradjah et al. 2015, Merazi et al. 2015, Al-Basyouni et al. 2015, Attia et al. 2015, Ait Atmane et al. 2015, Belkorissat et al. 2015, Boukhari et al. 2016, Bounouara et al. 2016, Mouaici et al. 2016, Beldjelili et al. 2016) proposed a new HSDTs to investigate the mechanical behavior of composite structures. Recently, Mantari and Granados (2015) proposed a new FSDT with four variables in which integral terms in the plate kinematics are utilized. Since FSDTs do not respect equilibrium conditions at the top and bottom surfaces of the plate, shear correction coefficients are needed to correct the unrealistic distribution of the shear strain/stress within the thickness. For these reasons, many HSDTs have been proposed to improve the limitations of FSDT such as Levinson (1980), Bhimaraddi and Stevens (1984), Reddy (1984), Ren (1986), Kant and Pandya (1988), and Mohan et al. (1994). A good review of these models for the investigation of laminated plates is found in (Noor and Burton 1989a, b, Reddy 1990 and 1993, Mallikarjuna and Kant 1993, Dahsin and Xiaoyu 1996). Reddy (1984) proposed a HSDT with cubic distributions for axial displacements. Based on Reddy's theory, Xiang et al. (2011) developed a n-order shear deformation theory. Kant and Pandya (1988), Mallikarjuna and Kant (1989) and Kant and Khare (1997) employed also HSDTs with cubic distributions for axial displacements as in the article by Reddy (1984). Recently, a new class of HSDTs is proposed by many researchers such as Shahrjerdi et al. (2011), Bouderba et al. (2013), Viswanathan et al. (2013), Ait Amar Meziane et al. (2014), Belabed et al. (2014), Ahmed (2014), Swaminathan and Naveenkumar (2014), Nedri et al. (2014), Bourada et al. (2015), Hamidi et al. (2015), Kar et al. (2015), Hebali et al. (2014), Mahi et al. (2015), Saidi et al. (2016), Bennoun et al. (2016) and Tounsi et al. (2016).

In the present research, the buckling analysis of isotropic and orthotropic plates under the inplane loading is investigated by utilizing a novel HSDT with four variables in which instead of derivative terms in the displacement field, integral terms are employed. The displacement field presented by Mantari and Granados (2015) is improved by considering higher-order distributions of in-plane displacements within the plate thickness. By employing the Navier procedure, the closed-form solutions have been determined. Numerical results considering side-to-thickness ratio and modulus ratio are carried out to illustrate the validity of the present model in computing the critical buckling load of isotropic and orthotropic plates. It can be concluded that the present model, which does not use shear correction coefficient, is not only simple but also comparable to the FSDT.

2. Theoretical formulation

2.1 Kinematics and strains

In this research work, further simplifying supposition are considered to the conventional HSDT so that the number of variables is diminished. The kinematic of the conventional HSDT is expressed by

$$u(x, y, z) = u_0(x, y) - z \frac{\partial w_0}{\partial x} + f(z)\varphi_x(x, y)$$
(1a)

$$v(x, y, z) = v_0(x, y) - z \frac{\partial w_0}{\partial y} + f(z)\varphi_y(x, y)$$
(1b)

$$w(x, y, z) = w_0(x, y)$$
 (1c)

where u_0 ; v_0 ; w_0 , φ_x , φ_y are five unknown displacements of the mid-plane of the plate, f(z) denotes shape function defining the distribution of the transverse shear strains and stresses within the thickness. By supposing that $\varphi_x = \int \theta(x, y) dx$ and $\varphi_y = \int \theta(x, y) dy$, the displacement field of the current theory model can be found in a simpler form as

$$u(x, y, z, t) = u_0(x, y, t) - z \frac{\partial w_0}{\partial x} + k_1 f(z) \int \theta(x, y, t) dx$$
(2a)

$$v(x, y, z, t) = v_0(x, y, t) - z \frac{\partial w_0}{\partial y} + k_2 f(z) \int \theta(x, y, t) dy$$
(2b)

$$w(x, y, z, t) = w_0(x, y, t)$$
 (2c)

In this investigation, the proposed higher-order shear deformation plate theory is determined by considering

$$f(z) = z \left(\frac{5}{4} - \frac{5z^2}{3h^2}\right)$$
(3)

It can be observed that the kinematic in Eq. (2) incorporates only four variables (u_0 , v_0 , w_0 and θ). The nonzero strains associated with the kinematic in Eq. (2) are

$$\begin{cases} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{cases} = \begin{cases} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \gamma_{xy}^0 \end{cases} + z \begin{cases} k_x^b \\ k_y^b \\ k_{xy}^b \end{cases} + f(z) \begin{cases} k_x^s \\ k_y^s \\ k_{xy}^s \end{cases}, \qquad \begin{cases} \gamma_{yz} \\ \gamma_{xz} \end{cases} = g(z) \begin{cases} \gamma_{yz}^0 \\ \gamma_{xz}^0 \end{cases}, \tag{4}$$

where

$$\begin{cases} \mathcal{E}_{x}^{0} \\ \mathcal{E}_{y}^{0} \\ \gamma_{xy}^{0} \end{cases} = \begin{cases} \frac{\partial u_{0}}{\partial x} \\ \frac{\partial v_{0}}{\partial x} \\ \frac{\partial u_{0}}{\partial y} + \frac{\partial v_{0}}{\partial x} \end{cases}, \quad \begin{cases} k_{x}^{b} \\ k_{y}^{b} \\ k_{xy}^{b} \end{cases} = \begin{cases} -\frac{\partial^{2} w_{0}}{\partial x^{2}} \\ -\frac{\partial^{2} w_{0}}{\partial y^{2}} \\ -2\frac{\partial^{2} w_{0}}{\partial x \partial y} \end{cases}, \quad (5a)$$

$$k_{x}^{s} \\ k_{y}^{s} \\ k_{xy}^{s} \end{cases} = \begin{cases} k_{1} \theta \\ k_{2} \theta \\ k_{1} \frac{\partial}{\partial y} \int \theta \ dx + k_{2} \frac{\partial}{\partial x} \int \theta \ dy \end{cases}, \quad \begin{cases} \gamma_{yz}^{0} \\ \gamma_{xz}^{0} \end{cases} = \begin{cases} k_{2} \int \theta \ dy \\ k_{1} \int \theta \ dx \end{cases},$$

and

$$g(z) = \frac{df(z)}{dz}$$
(5b)

The integrals used in the above expressions shall be resolved by a Navier type method and can be expressed as follows

$$\frac{\partial}{\partial y}\int\theta\ dx = A'\frac{\partial^2\theta}{\partial x\partial y}, \qquad \frac{\partial}{\partial x}\int\theta\ dy = B'\frac{\partial^2\theta}{\partial x\partial y}, \qquad \int\theta\ dx = A'\frac{\partial\theta}{\partial x}, \quad \int\theta\ dy = B'\frac{\partial\theta}{\partial y} \tag{6}$$

where the coefficients A' and B' are adopted according to the type of solution employed, in this case via Navier. Hence, A' and B' are defined as follows

$$A' = -\frac{1}{\alpha^2}, \quad B' = -\frac{1}{\beta^2}, \quad k_1 = \alpha^2, \quad k_2 = \beta^2$$
 (7)

where α and β are defined in expression (20).

2.2 Constitutive equations

For elastic and orthotropic plate, the constitutive relations can be expressed as

$$\begin{cases} \sigma_x \\ \sigma_y \\ \tau_{xy} \\ \tau_{yz} \\ \tau_{xz} \end{cases} = \begin{bmatrix} Q_{11} & Q_{12} & 0 & 0 & 0 \\ Q_{12} & Q_{22} & 0 & 0 & 0 \\ 0 & 0 & Q_{66} & 0 & 0 \\ 0 & 0 & 0 & Q_{44} & 0 \\ 0 & 0 & 0 & 0 & Q_{55} \end{bmatrix} \begin{bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \\ \gamma_{xy} \\ \gamma_{xz} \end{bmatrix}$$
(8)

where Q_{ij} present the plane stress reduced elastic constants in the material axes of the plate, and are expressed as

$$Q_{11} = \frac{E_1}{1 - v_{12}v_{21}}, \quad Q_{12} = \frac{v_{12}E_2}{1 - v_{12}v_{21}}, \quad Q_{22} = \frac{E_2}{1 - v_{12}v_{21}}, \quad Q_{66} = G_{12}, \quad Q_{44} = G_{23}, \quad Q_{55} = G_{13}$$
(9)

where E_1 , E_2 are Young's modulus, G_{12} , G_{23} , G_{13} are shear modulus, and v_{12} , v_{21} are Poisson's ratios. For the isotropic plate, these above material characteristics reduce to $E_1 = E_2 = E$, $G_{12} = G_{23} = G_{13} = G$, $v_{12} = v_{21} = v$. The subscripts 1, 2, 3 correspond to x, y, z directions of Cartesian coordinate system, respectively.

2.3 Governing equations

The principle of virtual works is employed herein to derive the governing equations. The principle can be expressed in analytical form as

$$\delta U + \delta V = 0 \tag{10}$$

where δU is the variation of strain energy; and δV is the variation of the external work done by external load applied to the plate.

The variation of strain energy of the plate is given by

$$\delta U = \int_{V} \left[\sigma_{x} \delta \varepsilon_{x} + \sigma_{y} \delta \varepsilon_{y} + \tau_{xy} \delta \gamma_{xy} + \tau_{yz} \delta \gamma_{yz} + \tau_{xz} \delta \gamma_{xz} \right] dV$$

$$= \int_{A} \left[N_{x} \delta \varepsilon_{x}^{0} + N_{y} \delta \varepsilon_{y}^{0} + N_{xy} \delta \gamma_{xy}^{0} + M_{x}^{b} \delta k_{x}^{b} + M_{y}^{b} \delta k_{y}^{b} + M_{xy}^{b} \delta k_{xy}^{b} + M_{xy}^{b} \delta k_{y}^{b} + M_{xy}^{b} \delta k_{xy}^{b} + M_{xy}^{b} \delta k_{y}^{b} + M_{xy}^{b} \delta k$$

where A is the top surface and the stress resultants N, M, and S are expressed by

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$$(N_i, M_i^b, M_i^s) = \int_{-h/2}^{h/2} (1, z, f) \sigma_i dz , \quad (i = x, y, xy) \text{ and } (S_{xz}^s, S_{yz}^s) = \int_{-h/2}^{h/2} g(\tau_{xz}, \tau_{yz}) dz$$
(12)

The variation of the external work can be expressed as

$$\delta V = -\int_{A} \left(N_x^0 \frac{\partial w_0}{\partial x} \frac{\partial \delta w_0}{\partial x} + 2N_{xy}^0 \frac{\partial w_0}{\partial x} \frac{\partial \delta w_0}{\partial y} + N_y^0 \frac{\partial w_0}{\partial y} \frac{\partial \delta w_0}{\partial y} \right) dA$$
(13)

where (N_x^0, N_y^0, N_{xy}^0) are transverse and in-plane applied loads, respectively. By substituting Eqs. (11) and (13) into Eq. (10), the following governing equations can be derived γ

$$\delta u_{0} : \frac{\partial N_{x}}{\partial x} + \frac{\partial N_{xy}}{\partial y} = 0$$

$$\delta v_{0} : \frac{\partial N_{xy}}{\partial x} + \frac{\partial N_{y}}{\partial y} = 0$$

$$\delta w_{0} : \frac{\partial^{2} M_{x}^{b}}{\partial x^{2}} + 2 \frac{\partial^{2} M_{xy}^{b}}{\partial x \partial y} + \frac{\partial^{2} M_{y}^{b}}{\partial y^{2}} + N_{x}^{0} \frac{\partial^{2} w_{0}}{\partial x^{2}} + 2 N_{xy}^{0} \frac{\partial^{2} w_{0}}{\partial x \partial y} + N_{y}^{0} \frac{\partial^{2} w_{0}}{\partial y^{2}} = 0$$

$$\delta \theta : -M_{x}^{s} - M_{y}^{s} - (A' + B') \frac{\partial^{2} M_{xy}^{s}}{\partial x \partial y} + A' \frac{\partial S_{xz}^{s}}{\partial x} + B' \frac{\partial S_{yz}^{s}}{\partial y} = 0$$
(14)

By substituting Eq. (4) into Eq. (8) and the subsequent results into Eq. (12), the stress resultants can be written as below

$$\begin{cases} N_x \\ N_y \\ N_{xy} \\ N_{xy} \\ M_x^b \\ M_y^b \\ M_x^b \\ M_y^b \\ M_x^s \\ M_y^s \\ M_x^s \\ M_y^s \\ M_x^s \\ M_y^s \\ M_x^s \\ M_x$$

$$\begin{cases} S_{xz}^{s} \\ S_{yz}^{s} \end{cases} = \begin{bmatrix} A_{55}^{s} & 0 \\ 0 & A_{44}^{s} \end{bmatrix} \begin{cases} \gamma_{xz}^{0} \\ \gamma_{yz}^{0} \end{cases}$$
(15b)

and stiffness components are expressed as

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$$\left\{A_{ij}, B_{ij}, D_{ij}, B_{ij}^{s}, D_{ij}^{s}, H_{ij}^{s}\right\} = \int_{-h/2}^{h/2} Q_{ij} \left\{1, z, z^{2}, f(z), z f(z), f^{2}(z)\right\} dz, \quad i, j = 1, 2, 6$$
(16a)

$$A_{ij}^{s} = \int_{-h/2}^{h/2} Q_{ij} [g(z)]^{2} dz, \quad i, j = 4,5$$
(16b)

Introducing relations (15) into Eq. (14), the governing equations can be expressed in terms of displacements (δu_0 , δv_0 , δw_0 , $\delta \theta$) and the appropriate equations take the form

$$A_{11}d_{11}u_0 + A_{66}d_{22}u_0 + (A_{12} + A_{66})d_{12}v_0 - (B_{11}d_{111}w_0 + (B_{12} + 2B_{66})d_{122}w_0) + (k_1A' + k_2B')B_{66}^s d_{122}\theta + (k_1B_{11}^s + k_2B_{12}^s)d_1\theta = 0,$$
(17a)

$$(A_{12} + A_{66})d_{12}u_0 + A_{66}d_{22}v_0 + A_{22}d_{22}v_0 - ((B_{12} + 2B_{66})d_{112}w_0 + B_{22}d_{222}w_0) + (k_1A' + k_2B')B_{66}^s d_{112}\theta + (k_2B_{22}^s + k_1B_{12}^s)d_2\theta = 0,$$
(17b)

$$(B_{11}d_{111}u_0 + (B_{12} + 2B_{66})d_{122}u_0) + ((B_{12} + 2B_{66})d_{112}v_0 + B_{22}d_{222}v_0) - D_{11}d_{1111}w_0 - 2(D_{12} + 2D_{66})d_{1122}w_0 - D_{22}d_{2222}w_0 + (k_1D_{11}^s + k_2D_{12}^s)d_{11}\theta + 2(k_1A' + k_2B')D_{66}^s d_{1122}\theta + (k_1D_{12}^s + k_2D_{22}^s)d_{22}\theta + N_x^0 d_{11}w_0 + 2N_{xy}^0 d_{12}w_0 + N_y^0 d_{22}w_0 = 0$$

$$(17c)$$

$$-\left(\left(k_{1}A'+k_{2}B'\right)B_{66}^{s}d_{122}u_{0}+\left(k_{1}B_{11}^{s}+k_{2}B_{12}^{s}\right)d_{1}u_{0}\right)-\left(\left(k_{1}A'+k_{2}B'\right)B_{66}^{s}d_{112}v_{0}+\left(k_{2}B_{22}^{s}+k_{1}B_{12}^{s}\right)d_{2}v_{0}\right) +\left(D_{11}^{s}+D_{12}^{s}\right)d_{11}w_{0}+2\left(k_{1}A'+k_{2}B'\right)D_{66}^{s}d_{1122}w_{0}+\left(D_{12}^{s}+D_{22}^{s}\right)d_{22}w_{0} -k_{1}^{2}H_{11}^{s}\theta-k_{2}^{2}H_{22}^{s}\theta-2k_{1}k_{2}H_{12}^{s}\theta-\left(k_{1}A'+k_{2}B'\right)^{2}H_{66}^{s}d_{1122}\theta +A_{54}^{s}\left(k_{2}B'\right)^{2}d_{22}\theta+A_{55}^{s}\left(k_{1}A'\right)^{2}d_{11}\theta=0$$
(17d)

where d_{ij} , d_{ijl} and d_{ijlm} are the following differential operators

$$d_{ij} = \frac{\partial^2}{\partial x_i \partial x_j}, \quad d_{ijl} = \frac{\partial^3}{\partial x_i \partial x_j \partial x_l}, \quad d_{ijlm} = \frac{\partial^4}{\partial x_i \partial x_j \partial x_l \partial x_m}, \quad d_i = \frac{\partial}{\partial x_i}, \quad (i, j, l, m = 1, 2).$$
(18)

2.4 Closed-form solution

The critical buckling loads of simply supported, orthotropic, rectangular plate will be computed in this work by employing the Navier method. Considering that the plate is subjected to in-plane compressive loads of form: $N_x^0 = \gamma_1 N_{cr}$, $N_y^0 = \gamma_2 N_{cr}$, $N_{xy}^0 = 0$ (here γ_1 and γ_2 are non-dimensional load parameters). The following displacement functions are chosen to automatically satisfy the boundary conditions

with

$$\alpha = m\pi/a \quad \text{and} \quad \beta = n\pi/b$$
 (20)

and U_{mn} , V_{mn} , W_{mn} , X_{mn} are arbitrary parameters to be determined. Substituting Eq. (19) into Eq. (17), the closed-form solution of buckling load can be determined from

$$\begin{bmatrix} S_{11} & S_{12} & S_{13} & S_{14} \\ S_{12} & S_{22} & S_{23} & S_{24} \\ S_{13} & S_{23} & S_{33} + k & S_{34} \\ S_{14} & S_{24} & S_{34} & S_{44} \end{bmatrix} \begin{bmatrix} U_{mn} \\ V_{mn} \\ W_{mn} \\ X_{mn} \end{bmatrix} = \begin{cases} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$
(21)

where

$$S_{11} = -\left(A_{11}\alpha^{2} + A_{66}\beta^{2}\right), \quad S_{12} = -\alpha\beta \left(A_{12} + A_{66}\right), \quad S_{13} = \alpha \left(B_{11}\alpha^{2} + B_{12}\beta^{2} + 2 B_{66}\beta^{2}\right), \\S_{14} = \alpha \left(k_{1}B_{11}^{s} + k_{2}B_{12}^{s} - (k_{1}A' + k_{2}B')B_{66}^{s}\beta^{2}\right), \\S_{22} = -\left(A_{66}\alpha^{2} + A_{22}\beta^{2}\right), \quad S_{23} = \beta \left(B_{22}\beta^{2} + B_{12}\alpha^{2} + 2B_{66}\alpha^{2}\right), \\S_{24} = \beta \left(k_{2}B_{22}^{s} + k_{1}B_{12}^{s} - (k_{1}A' + k_{2}B')B_{66}^{s}\alpha^{2}\right), \\S_{33} = -\left(D_{11}\alpha^{4} + 2(D_{12} + 2D_{66})\alpha^{2}\beta^{2} + D_{22}\beta^{4}\right), \\S_{34} = -\left(k_{1}D_{11}^{s} + k_{2}D_{12}^{s}\right)\alpha^{2} + 2\left(k_{1}A' + k_{2}B'\right)D_{66}^{s}\alpha^{2}\beta^{2} - \left(k_{2}D_{22}^{s} + k_{1}D_{12}^{s}\right)\beta^{2}, \\S_{44} = -k_{1}\left(k_{1}H_{11}^{s} + k_{2}H_{12}^{s}\right) - \left(k_{1}A' + k_{2}B'\right)^{2}\left(H_{66}^{s}\alpha^{2}\beta^{2}\right) - k_{2}\left(k_{1}H_{12}^{s} + k_{2}H_{22}^{s}\right) \\- \left(k_{1}A'\right)^{2}A_{55}^{s}\alpha^{2} - \left(k_{2}B'\right)^{2}A_{44}^{s}\beta^{2} \\k = N_{cr}\left(\gamma_{1}\alpha^{2} + \gamma_{2}\beta^{2}\right)$$

$$(22)$$

The geometrical instability has been taken care in the present analysis by considering the determinant of the coefficient matrix in Eq. (21) to be equal zero. Indeed, for nontrivial solution, the determinant of the coefficient matrix in Eq. (21) must be zero. For each choice of m and n, there is a corresponding unique value of N_{cr} . The critical buckling load is the smallest value of N_{cr} (m, n).

By applying the condensation approach to eliminate the in-plane displacements U_{mn} and V_{mn} , Eq. (21) can be rewritten as

$$\begin{bmatrix} \overline{S}_{33} + k & \overline{S}_{34} \\ \overline{S}_{43} & \overline{S}_{44} \end{bmatrix} \begin{bmatrix} W_{mn} \\ X_{mn} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
(23)

where

$$\overline{S}_{33} = S_{33} - \frac{S_{13}(S_{13}S_{22} - S_{12}S_{23}) - S_{23}(S_{11}S_{23} - S_{12}S_{13})}{S_{11}S_{22} - S_{12}^2}$$

$$\overline{S}_{34} = S_{34} - \frac{S_{14}(S_{13}S_{22} - S_{12}S_{23}) - S_{24}(S_{11}S_{23} - S_{12}S_{13})}{S_{11}S_{22} - S_{12}^2}$$

$$\overline{S}_{43} = S_{34} - \frac{S_{13}(S_{14}S_{22} - S_{12}S_{24}) - S_{23}(S_{11}S_{24} - S_{12}S_{14})}{S_{11}S_{22} - S_{12}^2}$$

$$\overline{S}_{44} = S_{44} - \frac{S_{14}(S_{14}S_{22} - S_{12}S_{24}) - S_{24}(S_{11}S_{24} - S_{12}S_{14})}{S_{11}S_{22} - S_{12}^2}$$

$$(24)$$

The system of homogeneous Eq. (23) has a nontrivial solution only for discrete values of the buckling load. For a nontrivial solution, the determinant of the coefficients (W_{mn}, X_{mn}) must equal zero

$$\begin{vmatrix} \overline{S}_{33} + k & \overline{S}_{34} \\ \overline{S}_{43} & \overline{S}_{44} \end{vmatrix} = 0$$
(25)

The resulting equation may be solved for the buckling load. This gives the following expression for buckling load

$$k = \frac{\overline{S}_{34}\overline{S}_{43} - \overline{S}_{33}\overline{S}_{44}}{\overline{S}_{44}}$$
(26)

By employing the Eq. (25), the following expression for critical buckling load is determined

$$N_{cr} = \frac{1}{\left(\gamma_1 \,\alpha^2 + \gamma_2 \beta^2\right)} \frac{\overline{S}_{34} \,\overline{S}_{43} - \overline{S}_{33} \,\overline{S}_{44}}{\overline{S}_{44}} \tag{27}$$

3. Numerical results and discussions

In this section the accuracy of the proposed HSDT which uses a kinematic with only four variables, is assessed. For this end, a simply supported rectangular plate under loading conditions, as presented in Fig. 1, is examined to demonstrate the accuracy of the current formulation in investigating the buckling response of the plate. For proposed examples, the following engineering constants are employed (Reddy 1997)





 E_1 / E_2 varied $G_{12} / E_2 = G_{13} / E_2 = 0.5$, $G_{23} / E_2 = 0.2$, $v_{12} = 0.25$ (28)

For convenience, the following non-dimensional buckling load is utilized

$$\overline{N} = \frac{N_{cr}a^2}{E_2h^3} \tag{29}$$

where *a* is the length of the square plate and *h* is the thickness of the plate.

a / h	Theories	Isotropic $v = 0.3$	Orthotropic		
			$E_1/E_2 = 10$	$E_1 / E_2 = 25$	$E_1 / E_2 = 40$
	Present	2.9512	6.3478	9.1039	10.5785
	RPT ^(*)	2.9512	6.3478	9.1039	10.5785
5	$\mathrm{FSDT}^{(*)} \left(k_s = 2 / 3 \right)$	2.8200	5.5679	7.1122	7.7411
	$\mathrm{FSDT}^{(*)} \left(k_s = 5 / 6 \right)$	2.9498	6.1804	8.2199	9.1085
	$FSDT^{(*)}(k_s=1)$	3.0432	6.6715	9.1841	10.3463
	Present	3.4224	9.3732	16.7719	22.2581
	RPT ^(*)	3.4224	9.3732	16.7719	22.2581
10	$\mathrm{FSDT}^{(*)}(k_s=2/3)$	3.3772	8.8988	14.7011	18.3575
	$\mathrm{FSDT}^{(*)}(k_s=5/6)$	3.4222	9.2733	15.8736	20.3044
	$FSDT^{(*)}(k_s=1)$	3.4530	9.5415	16.7699	21.8602
20	Present	3.5650	10.6534	21.3479	31.0685
	RPT ^(*)	3.5650	10.6534	21.3479	31.0685
	$\mathrm{FSDT}^{(*)}(k_s=2/3)$	3.5526	10.4926	20.4034	28.85
	$\mathrm{FSDT}^{(*)}\left(k_{s}=5/6\right)$	3.5650	10.6199	20.9528	30.0139
	$FSDT^{(*)}(k_s=1)$	3.5733	10.7066	21.3363	30.8451

Table 1 Comparison of non-dimensional critical buckling load of square plates under uniaxial compression

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	Present	3.6071	11.0780	23.1225	34.9717
	RPT ^(*)	3.6071	11.0780	23.1225	34.9717
50	$FSDT^{(*)}(k_s = 2/3)$	3.6051	11.0497	22.9366	34.4886
	$FSDT^{(*)}(k_s = 5/6)$	3.6071	11.0721	23.0461	34.7487
	$FSDT^{(*)}(k_s=1)$	3.6085	11.0871	23.1197	34.9244
	Present	3.6132	11.1415	23.4007	35.6120
100	RPT ^(*)	3.6132	11.1415	23.4007	35.6120
	$FSDT^{(*)}(k_s = 2/3)$	3.6127	11.1343	23.3527	35.4852
	$FSDT^{(*)}(k_s = 5/6)$	3.6132	11.1400	23.3810	35.5538
	$FSDT^{(*)}(k_s = 1)$	3.6135	11.1438	23.3999	35.5996

* Taken from Ref (Kim *et al.* 2009) Table 2 Comparison of non-dimensional critical buckling load of square plates under biaxial compression

a / h	Theories	Isotropic	Orthotropic		
		v = 0.3	$E_1 / E_2 = 10$	$E_1 / E_2 = 25$	$E_1 / E_2 = 40$
	Present	1.4756	2.8549 ^a	3.3309 ^a	3.3800 ^a
	RPT ^(*)	1.4756	2.8549^{a}	3.3309 ^a	3.3800 ^a
5	$FSDT^{(*)}(k_s = 2/3)$	1.4100	2.5042^{a}	2.7332 ^a	2.8303 ^a
	$FSDT^{(*)}(k_s = 5/6)$	1.4749	2.8319 ^a	3.1422 ^a	3.2822 ^a
	$FSDT^{(*)}(k_s=1)$	1.5216	3.1027 ^a	3.4933 ^a	3.6793 ^a
	Present	1.7112	4.6718 ^a	6.0646 ^a	7.2536 ^a
	RPT ^(*)	1.7112	4.6718 ^a	6.0646 ^a	7.2536 ^a
10	$\mathrm{FSDT}^{(*)}(k_s=2/3)$	1.6886	4.4259	5.4351 ^a	6.0797 ^a
	$FSDT^{(*)}(k_s = 5/6)$	1.7111	4.6367	5.8370 ^a	6.6325 ^a
	$FSDT^{(*)}(k_s=1)$	1.7265	4.7708	6.1425 ^a	7.0690^{a}
	Present	1.7825	5.3267	7.6643 ^a	9.6614 ^a
	RPT ^(*)	1.7825	5.3267	7.6643 ^a	9.6614 ^a
20	$FSDT^{(*)}(k_s = 2/3)$	1.7763	5.2463	7.3701 ^a	8.9895ª
	$FSDT^{(*)}(k_s = 5/6)$	1.7825	5.3100	7.5546 ^a	9.3049 ^a
	$\mathrm{FSDT}^{(*)}(k_s=1)$	1.7866	5.3533	7.6834 ^a	9.5297 ^a
	Present	1.8036	5.5390	8.2784 ^a	10.6576 ^a
	RPT ^(*)	1.8036	5.5390	8.2784 ^a	10.6576 ^a
50	$\mathrm{FSDT}^{(*)}(k_s=2/3)$	1.8025	5.5249	8.2199 ^a	10.5111 ^a
	$FSDT^{(*)}(k_s = 5/6)$	1.8036	5.5361	8.2566 ^a	10.5810^{a}
	$FSDT^{(*)}(k_s=1)$	1.8042	5.5436	8.2812 ^a	10.6282 ^a
100	Present	1.8066	5.5707	8.3744 ^a	10.8172 ^a
	RPT ^(*)	1.8066	5.5707	8.3744 ^a	10.8172 ^a
	$\mathrm{FSDT}^{(*)}\left(k_{s}=2/3\right)$	1.8063	5.5672	8.3593 ^a	10.7788^{a}
	$FSDT^{(*)}(k_s = 5/6)$	1.8066	5.5700	8.3687 ^a	10.7972 ^a
	$FSDT^{(*)}(k_s = 1)$	1.8068	5.5719	8.3751 ^a	10.8095 ^a

* Taken from Ref (Kim et al. 2009)

^a Mode for plate is (m, n) = (1, 2).

Example 1: Table 1 gives the values of the non-dimensional buckling loads \overline{N} of isotropic and orthotropic square plates subjected to uniaxial compression for various values of thickness ratio a/h and modulus ratio E_1/E_2 . The obtained results are compared with the data reported by Kim *et al.* (2009) based on FSDT and refined plate theory (RPT). Many shear correction coefficients ($k_s = 2/3$, $k_s = 5/6$ and $k_s = 1$) are employed for the FSDT (Kim *et al.* 2009) in comparison with the present model. An excellent agreement is proved for all types ranging from thin to very thick plates.

Example 2: In Table 2 we found the values of the non-dimensional buckling loads \overline{N} of isotropic and orthotropic square plates under biaxial compression for various values of thickness ratio a / h and modulus ratio E_1/E_2 . Again, the computed values are compared with those given by Kim *et al.* (2009) based on FSDT and RPT. It is remarked that there is an excellent agreement between the values computed by the present model and RPT (Kim *et al.* 2009) for all values of thickness ratio a / h and modulus ratio E_1/E_2 .

Example 3: This example is performed for isotropic and orthotropic square plates subjected to tension in the x direction and compression in the y direction. In Table 3 the comparison of nondimensional buckling loads \overline{N} computed via present model with those provided by Kim *et al.*

a / h	Theories	Isotropic $v = 0.3$	Orthotropic		
			$E_1/E_2 = 10$	$E_1 / E_2 = 25$	$E_1 / E_2 = 40$
	Present	4.8274 ^a	4.0258 ^b	4.1044 ^c	4.1525 ^c
	RPT ^(*)	4.8274 ^a	4.0258 ^b	4.1044 ^c	4.1525 ^c
5	$\mathrm{FSDT}^{(*)}\left(k_{s}=2/3\right)$	4.4175 ^a	3.2849 ^d	3.3001 ^e	3.3053 ^e
	$\mathrm{FSDT}^{(*)}\left(k_{s}=5/6\right)$	4.8158 ^a	3.9241 ^c	3.9794 ^c	4.0075 ^d
	$FSDT^{(*)}(k_s=1)$	5.1237 ^a	4.4488^{b}	4.5691 ^c	4.6073 ^c
	Present	6.6024 ^a	7.7863	8.5471 ^b	9.1638 ^b
	RPT ^(*)	6.6024 ^a	7.7863	8.5471 ^b	9.1638 ^b
10	$\mathrm{FSDT}^{(*)}\left(k_{s}=2/3\right)$	6.4032 ^a	7.2656	7.7820 ^b	8.1208 ^b
	$\mathrm{FSDT}^{(*)}\left(k_{s}=5/6\right)$	6.6010 ^a	7.7748	8.4774 ^b	8.9039 ^b
	$FSDT^{(*)}(k_s=1)$	6.7398 ^a	8.0651	9.0153 ^b	9.5197 ^b
20	Present	7.2754 ^a	9.2811	11.6347 ^b	12.8031 ^b
	RPT ^(*)	7.2754 ^a	9.2811	11.6347 ^b	12.8031 ^b
	$\mathrm{FSDT}^{(*)}(k_s=2/3)$	7.2139 ^a	9.1310	11.2544 ^b	12.1990 ^b
	$\mathrm{FSDT}^{(*)}\left(k_{s}=5/6\right)$	7.2753 ^a	9.2782	11.6015 ^b	12.6339 ^b
	$FSDT^{(*)}(k_s=1)$	7.3168 ^a	9.3790	11.8453 ^b	12.9428 ^b
50	Present	7.4895 ^a	9.8101	12.9531 ^b	14.4177 ^b
	RPT ^(*)	7.4895 ^a	9.8101	12.9531 ^b	14.4177 ^b

Table 3 Comparison of non-dimensional critical buckling load of square plates subjected to tension in the *x* direction and compression in the *y* direction

	$FSDT^{(*)}(k_s = 2/3)$	7.4790 ^a	9.7830	12.8751 ^b	14.2839 ^b
	$FSDT^{(*)}(k_s = 5/6)$	7.4895^{a}	9.8097	12.9463 ^b	14.3789 ^b
	$FSDT^{(*)}(k_s=1)$	7.4965 ^a	9.8275	12.9942 ^b	14.4430 ^b
	Present	7.5211 ^a	9.8907 ^a	13.1666 ^b	14.6827 ^b
	RPT ^(*)	7.5211 ^a	9.8907 ^a	13.1666 ^b	14.6827 ^b
100	$FSDT^{(*)}(k_s = 2/3)$	7.5185 ^a	9.8838 ^a	16.1463 ^b	14.6474 ^b
	$FSDT^{(*)}(k_s = 5/6)$	7.5211 ^a	9.8906 ^a	13.1648 ^b	14.6724 ^b
	$FSDT^{(*)}(k_s=1)$	7.5229 ^a	9.8951 ^a	13.1772 ^b	14.6891 ^b

* Taken from Ref (Kim et al. 2009)

^a Mode for plate is (m, n) = (1, 2); ^b Mode for plate is (m, n) = (1, 3); ^c Mode for plate is (m, n) = (1, 4)^d Mode for plate is (m, n) = (1, 5); ^e Mode for plate is (m, n) = (1, 6).

(2009) based on FSDT and RPT is presented. It can be seen that there is an excellent agreement for a wide range of values of thickness ratio and modulus ratio.

Fig. 2 demonstrates the effect of thickness and modulus ratios on the non-dimensional critical buckling load N of square plate subjected to uniaxial compression. It is observed that the present novel four variable refined plate theory and RPT predict almost the same data, and CPT overestimates the buckling loads of plate due to neglecting transverse shear deformation influences. The difference between CPT and shear deformation models diminishes when the side-to-thickness ratio a / h increases. As well as the plate becomes more orthotropic, the difference between the present theory and FSDT will increase with respect to the increase of thickness ratio.

The variations of non-dimensional critical buckling load N of square plate as a function of the modulus ratio is illustrated in Fig. 3. The considered plate is subjected to biaxial compression. The curves plotted by employing the present model are compared with the curves plotted by utilizing the CPT, FSDT and the RPT. From these results can be concluded that the resulting curves are very close to the curves plotted by using RPT and the CPT over-estimates the results when the





Fig. 2 The effectofside-to thickness and modulus ratios on the critical buckling load of square plate subjected to uniaxial compression: (a) isotropic; (b) $E_1 / E_2 = 10$; (c) $E_1 / E_2 = 25$; and (d) $E_1 / E_2 = 40$



Fig. 3 The effect of modulus ratio on the critical buckling load of square plate subjected to uniaxial compression: (a) a = 10 h; and (b) a = 20 h



Fig. 4 The effect of modulus ratio on the critical buckling load of square plate subjected to biaxial compression: (a) a = 10 h ; and (b) a = 20 h

modulus ratio increases. It can be also seen that the difference between the present theory and SDT will increase with increasing the modulus ratio.

The effect of modulus ratio on the critical buckling load of square plate subjected to tension in the x direction and compression in the y direction is shown in Fig. 4. Again from Fig. 4 can be observed that the resulting curves are very close to the curves plotted by using RPT and the CPT over-estimates the results when the modulus ratio increases.

It can be observed from Tables 1-3 and Figs. 2-4 that the difference of critical buckling load between the present theory and FSDT depends on not only the thickness and modulus ratios, but also the in-plane loading conditions (Fig. 1).

4. Conclusions

A simplified HSDT with only four unknowns was developed for buckling analysis of isotropic and orthotropic plates. Governing equations are obtained from the principle of virtual works. Closed-form solutions are obtained for simply supported orthotropic plates. The accuracy of the developed model has been checked for the stability analysis of isotropic and orthotropic plates.

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